

Variational Quantum Eigensolver and Quantum Approximate Optimization Algorithm

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- quantum mechanics has unique features (superposition of states, entanglement, ...)
- exploit these features to perform computations
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 - adiabatic quantum computing and quantum annealing (analog approaches)
 - quantum cellular automata
 - ...

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- limited usefulness due to hardware constraints
- hybrid approaches:
 - variational quantum eigensolver and quantum approximate optimization algorithm
 - “robust” against gate imperfections, shallow circuits, already useful for a limited number of qubits

Notation

- notation for the Pauli matrices

$$I = \sigma^0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$Y = \sigma^2 = \sigma^y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$X = \sigma^1 = \sigma^x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad (1)$$

$$Z = \sigma^3 = \sigma^z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad (2)$$

- subscripts indicate qubit: $X_1 Y_2 = X \otimes Y$
- computational basis is z-basis:

$$|0\rangle = |\downarrow\rangle \quad (3)$$

$$|1\rangle = |\uparrow\rangle \quad (4)$$

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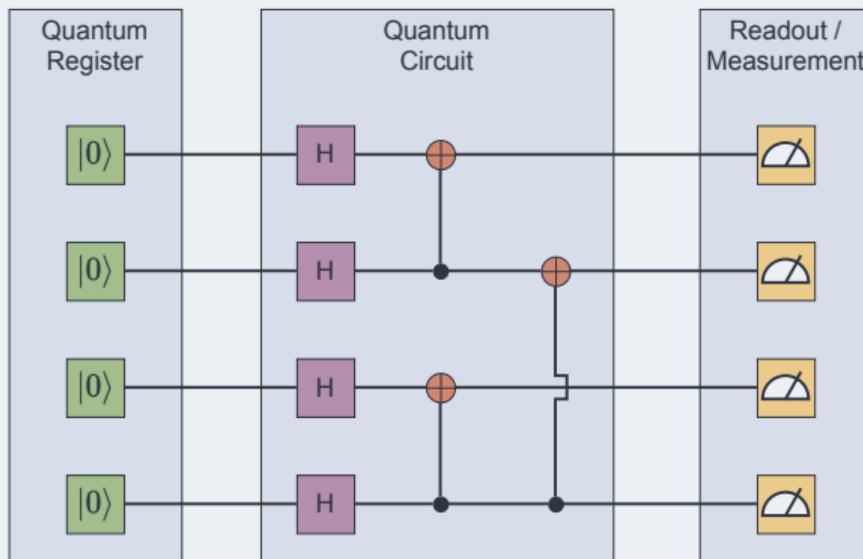
Max-Cut

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Digital Quantum Computer

Digital QC: Principle

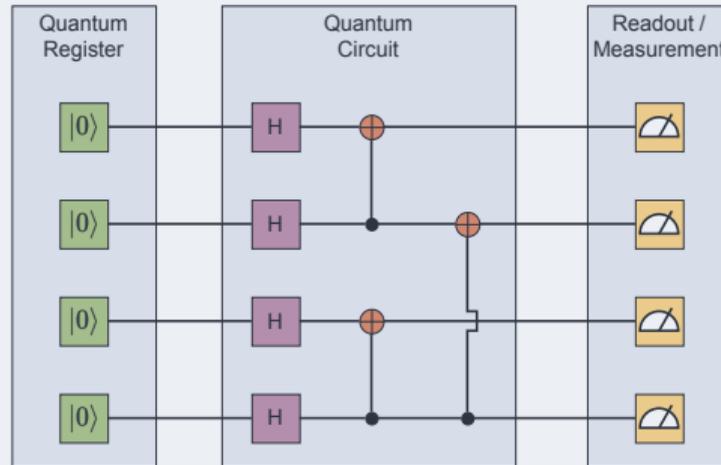
- applies quantum gates to registers of qubits, performs arbitrary computations
- interesting algorithms (e.g. Shor or Grover) have been proposed to solve problems faster than on a classical computer



Digital QC: Vincenzo's Criteria

Criteria for a universal quantum computer

- well-defined qubits (two-level system)
- initialization of the qubits in a well defined-state
- long decoherence time
- a universal set of quantum gates (very small set of one- and two-qubit gates)
- qubit-selective readout



Digital QC: Qubit Implementations

- many platforms have been suggested/explored:
 - Josephson junctions
 - photonic qubits
 - Rydberg atoms
 - electrons
 - nuclear spin (NMR)
 - quantum dots
 - topological systems (non-abelian anyons)
 - optical lattices (internal atomic states)
 - van der Waals heterostructures
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 - ...
- platforms differ vastly regarding the relevant characteristics
 - gate fidelities
 - connectivity
 - storage and decoherence time
 - scalability (number of qubits)

Digital QC: Characteristics

Advantages:

- can run arbitrary quantum algorithms (quantum Turing machine)

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- number of proposed algorithms is fairly limited
- so far only one claim of actual quantum advantage:
 - Google 2019: sampling random quantum circuits
 - 53 superconducting qubits
 - classical supercomputer was limited to 43 qubits

Analog Quantum Computer

Analog QC: Principle I

Here: Analog Quantum Computing = Adiabatic Quantum Computing

- solve an optimization problem encoded in a problem Hamiltonian H_P
- start with the ground state of a well known Hamiltonian

$$H_0 = \sum_{i=1}^N X_i \quad (5)$$

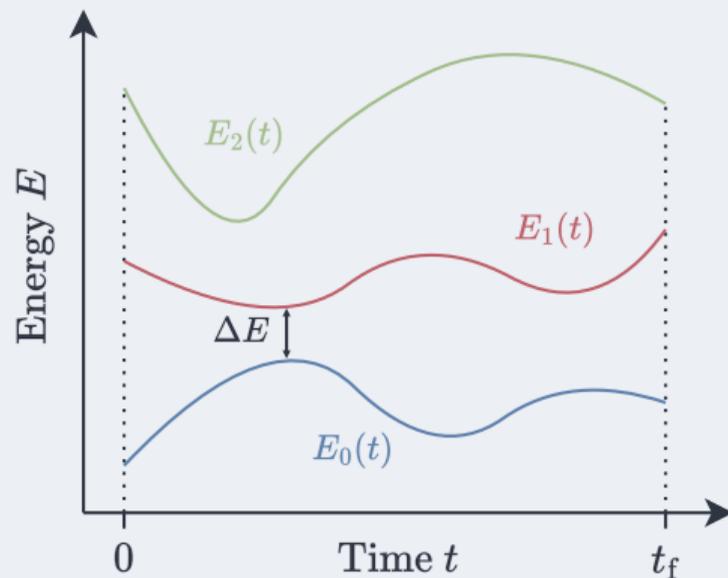
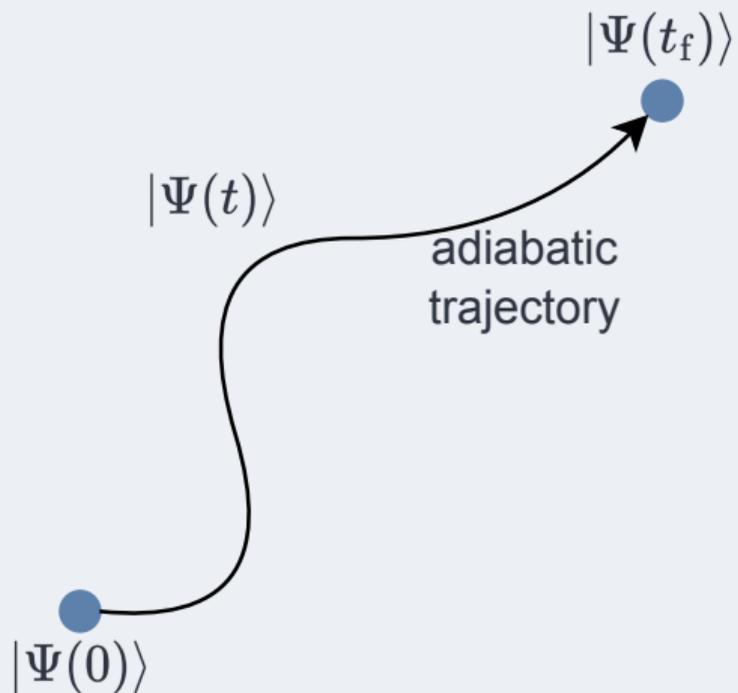
- slowly drive the Hamiltonian to H_P

$$H(t) = (1 - f(t))H_0 + f(t)H_P, \quad f(0) = 0, \quad f(t_f) = 1 \quad (6)$$

- if done slow enough the system will remain in the instantaneous ground state (adiabatic theorem)
- maximum feasible speed depends on minimum energy gap in the system

$$t_f = \mathcal{O}(\Delta E^{-2}) \quad (7)$$

Analog QC: Principle II



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- well suited for optimization/eigenvalue problems
- commercial platforms with thousands of qubits available (e.g. D-Wave with superconducting qubits)

Disadvantages:

- not universal, at least on current hardware (e.g. Shor algorithm cannot be mapped)
- near degeneracies/avoided crossings require slow driving
- no evidence for quantum speed-up so far

Variational Quantum Eigensolver (VQE)

VQE: Principle

- hybrid approach that may render today's quantum hardware more useful
- solve optimization/eigenvalue problems
- combine a quantum computer with a classical optimizer

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- variational principle guarantees

$$\langle \Psi(\vec{\theta}) | H_P | \Psi(\vec{\theta}) \rangle \geq E_0 \quad (8)$$

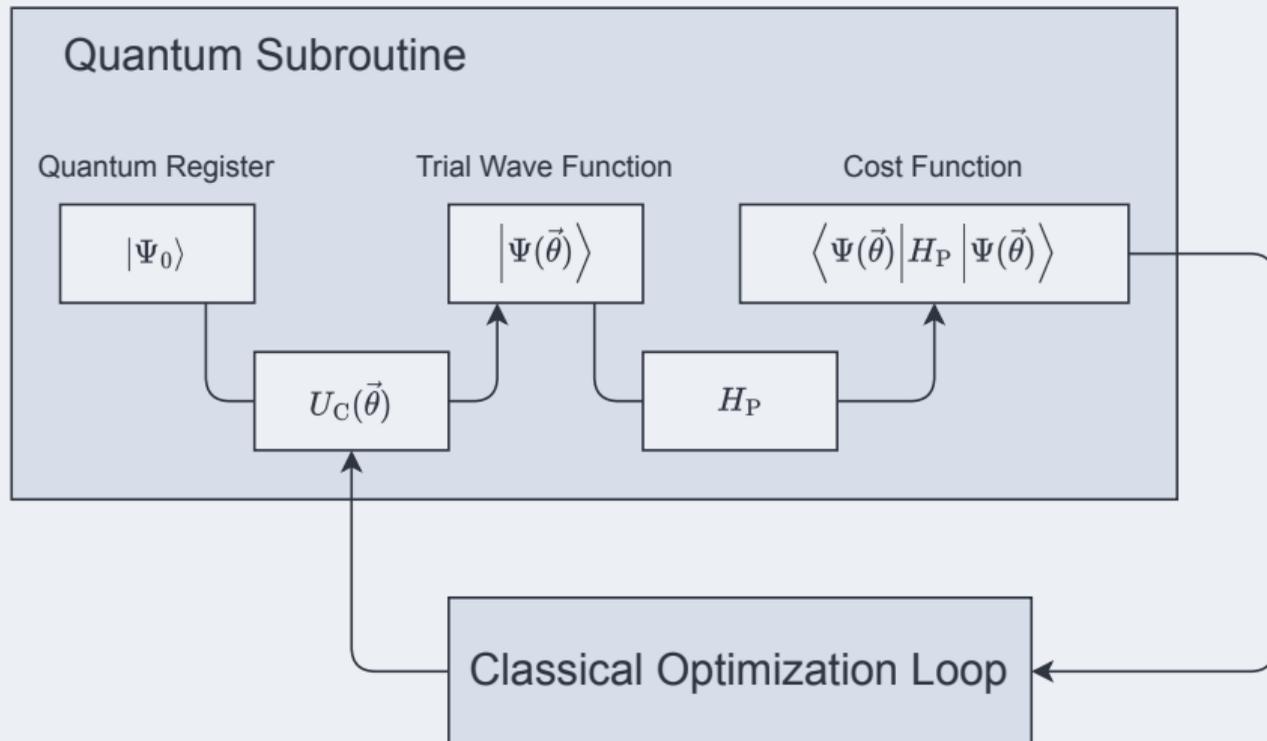
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- the optimization of parameters can mitigate imperfections of the hardware to some extent

VQE: Feedback Loop



VQE: Typical Structure of a Parametric Circuit

N qubits store the trial states

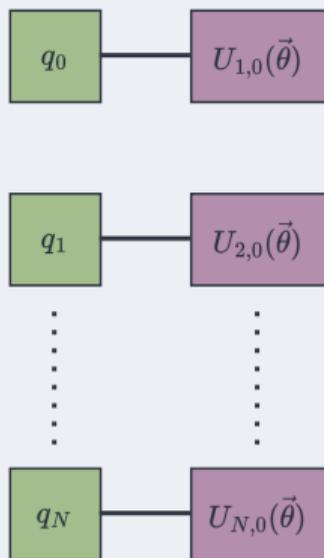


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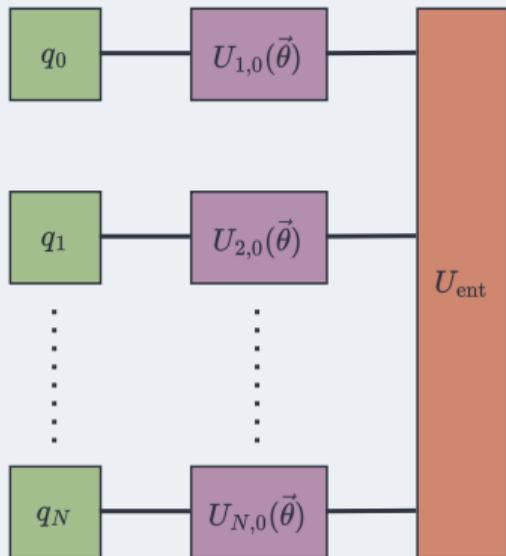
VQE: Typical Structure of a Parametric Circuit

single qubit gates rotate each qubit according to some parameters $\vec{\theta}$



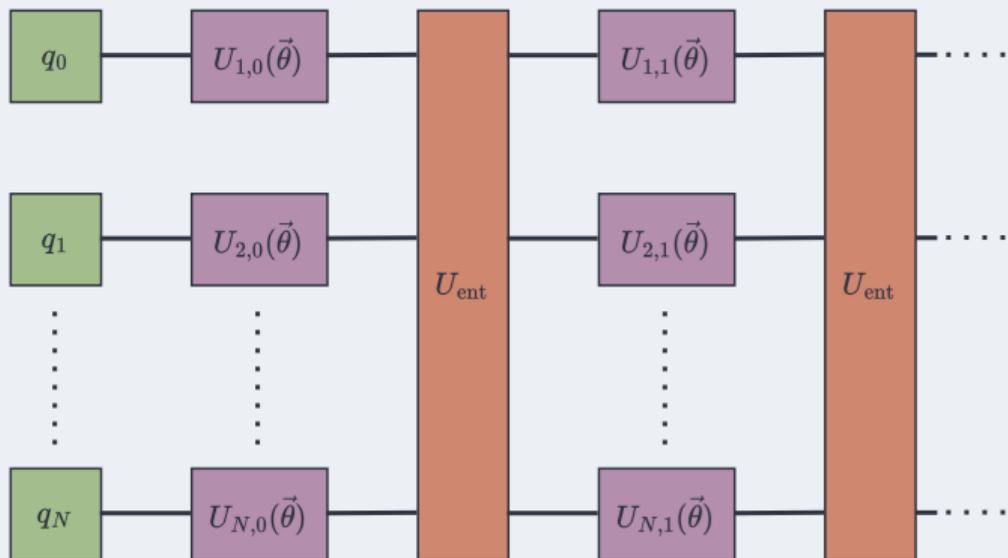
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entanglers generate entanglement among the qubits, crucial for quantum speed-up



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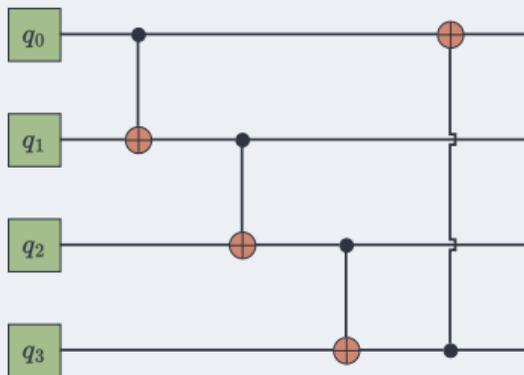
operator structure is repeated



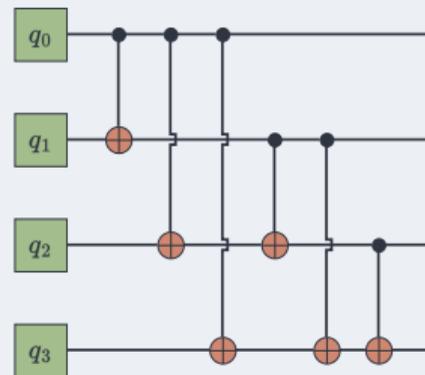
VQE: Entanglement Schemes



linear



circular



full

- of course other schemes are conceivable as well
- best choice depends on the connectivity of the hardware platform and the problem under consideration
- it would be nice to have a simple scheme guiding the construction of such circuits

VQE: Molecular Hamiltonian

- electronic problem in Born-Oppenheimer approximation

$$H = \underbrace{-\sum_{i=1}^{N_e} \frac{\hbar^2}{2m_e} \nabla^2}_{\text{kinetic}} - \underbrace{\sum_{i=1}^{N_n} \sum_{j=1}^{N_e} \frac{Z_i e^2}{4\pi\epsilon_0 |\vec{R}_i - \vec{r}_j|}}_{\text{electron-nuclei int.}} + \underbrace{\frac{1}{2} \sum_{\substack{i,j=1 \\ i \neq j}}^{N_e} \frac{e^2}{4\pi\epsilon_0 |\vec{r}_i - \vec{r}_j|}}_{\text{electron-electron int.}} \quad (9)$$

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- limited usefulness, does not incorporate electronic correlations
- can serve as a starting point for post-Hartree-Fock methods like configuration interaction or coupled cluster
- these methods are computationally hard due to the scaling w.r.t. system size
- only feasible for “small” molecules, typically up to $\mathcal{O}(10^2)$ electrons

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- quantum computers can efficiently store and work with these correlated states

VQE: Unitary Coupled Cluster (UCC) Ansatz I

- wave function ansatz

$$|\Psi\rangle = e^{T-T^\dagger} |\Psi_{\text{ref}}\rangle \quad (10)$$

- reference state $|\Psi_{\text{ref}}\rangle$, typically a single Slater determinant of Hartree-Fock orbitals
- cluster operator

$$T = \sum_{i=1}^{N_e} T_i \quad (11)$$

- terms describe i -particle excitations to unoccupied (virtual) orbitals

$$T_1 = \sum_{\substack{i \in \text{occ.} \\ a \in \text{virt.}}} t_a^i a_a^\dagger a_i \quad (12)$$

$$T_2 = \sum_{\substack{i > j \in \text{occ.} \\ a > b \in \text{virt.}}} t_{ab}^{ij} a_a^\dagger a_b^\dagger a_i a_j \quad (13)$$

$$\dots \quad (14)$$

VQE: Unitary Coupled Cluster Ansatz (UCC) II

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$$H = \sum_{pq} h_{pq} a_p^\dagger a_q + \sum_{pqkl} h_{pqkl} a_p^\dagger a_q^\dagger a_k a_l \quad (15)$$

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How to map the fermionic UCC function and the molecular Hamiltonian to a quantum computer (spin-1/2)?

VQE: Fermionic Transformations

- the fermionic operators a_j^\dagger and a_j can be transformed to spin-1/2 operators using the Jordan-Wigner transformation

$$a_j^\dagger = \exp \left[i\pi \sum_{k=1}^{j-1} \sigma_k^+ \sigma_k^- \right] \sigma_j^+ \quad (16)$$

$$a_j = \exp \left[-i\pi \sum_{k=1}^{j-1} \sigma_k^+ \sigma_k^- \right] \sigma_j^- \quad (17)$$

with $\sigma_j^+ = 1/2(X_j + iY_j)$ and $\sigma_j^- = 1/2(X_j - iY_j)$

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- alternative: Bravyi-Kitaev (BK) transformation
- example: BK trafo for \mathbb{H}_2

$$H = g_0 1 + g_1 Z_0 + g_2 Z_1 + g_3 Z_0 Z_1 + g_4 Y_0 Y_1 + g_5 X_0 X_1 \quad (18)$$

(parameters g_i depend on nuclear coordinates and some constants)

VQE: Quantum Expectation Estimation I

How do we measure $\langle \Psi(\vec{\theta}) | H_P | \Psi(\vec{\theta}) \rangle$?

- a quantum computer must be able to selectively measure Z_i
- other operators can be measured by performing unitary transformations U

$$O = U^\dagger Z U \quad (19)$$

- this allows for single-qubit Pauli measurements (measure X, Y, Z)

$$O = X \quad U = H \quad (20)$$

$$O = Y \quad U = HS^\dagger \quad (21)$$

$$O = Z \quad U = 1 \quad (22)$$

here H denotes the Hadamard and S the $\pi/4$ phase gate

- the same principle can be applied to multi-qubit measurements

$$O = X \otimes Z \qquad U = \text{CNOT} (H \otimes 1) \qquad (23)$$

$$O = X \otimes Y \qquad U = \text{CNOT} (H \otimes HS^\dagger) \qquad (24)$$

\vdots

\vdots

- this allows general Pauli strings to be measured, e.g. $X \otimes Y \otimes Y \otimes Z$

VQE: Quantum Expectation Estimation III

- let's assume the Hamiltonian consists of M Pauli strings

$$H = \sum_{\gamma=1}^M H_{\gamma} \quad (25)$$

- it's expectation value is given by

$$\langle H \rangle = \sum_{\gamma=1}^M \langle H_{\gamma} \rangle \quad (26)$$

- individual state preparation for every Pauli string \Rightarrow additive variances

$$\text{Var}[H] = \sum_{\gamma=1}^M \text{Var}[H_{\gamma}] \quad (27)$$

VQE: Quantum Expectation Estimation IV

- assuming a normal distribution

$$n_{\text{preparations}} = M \sum_{\gamma=1}^M \frac{\text{Var}[H_{\gamma}]}{\varepsilon^2} \quad (28)$$

state preparations are required to achieve a variance smaller than ε

- usually the $\text{Var}[H_{\gamma}]$ are unknown (but can be estimated)
- various strategies can be used to reduce the number of state preparations
- commuting terms can be measured in parallel
- for best results only terms with $\text{Cov}[H_{\alpha}, H_{\beta}] = 0$ should be grouped, otherwise correlations might increase the number of necessary measurements

$$\text{Example : } H = -X_0X_1 - Y_0Y_1 + Z_0Z_1 + Z_0 + Z_1 \quad (29)$$

$$\Rightarrow \{-X_0X_1\}, \{-Y_0Y_1, Z_0Z_1\}, \{Z_0, Z_1\} \quad (30)$$

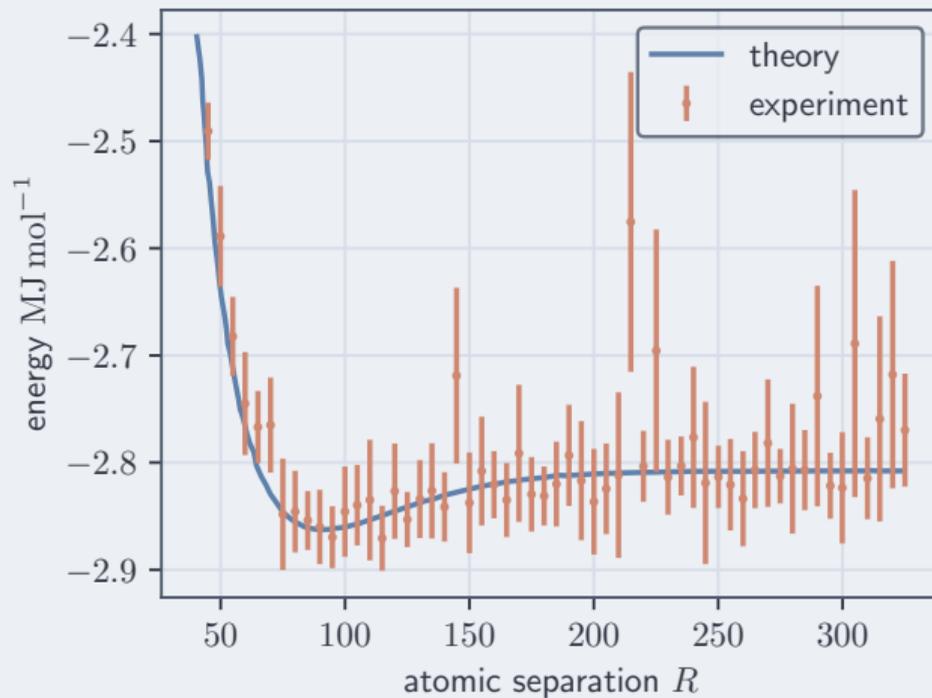
(covariance between X_0X_1 and Y_0Y_1)

VQE: Ingredients for Molecular Problem

We now have all ingredients for the electronic structure problem:

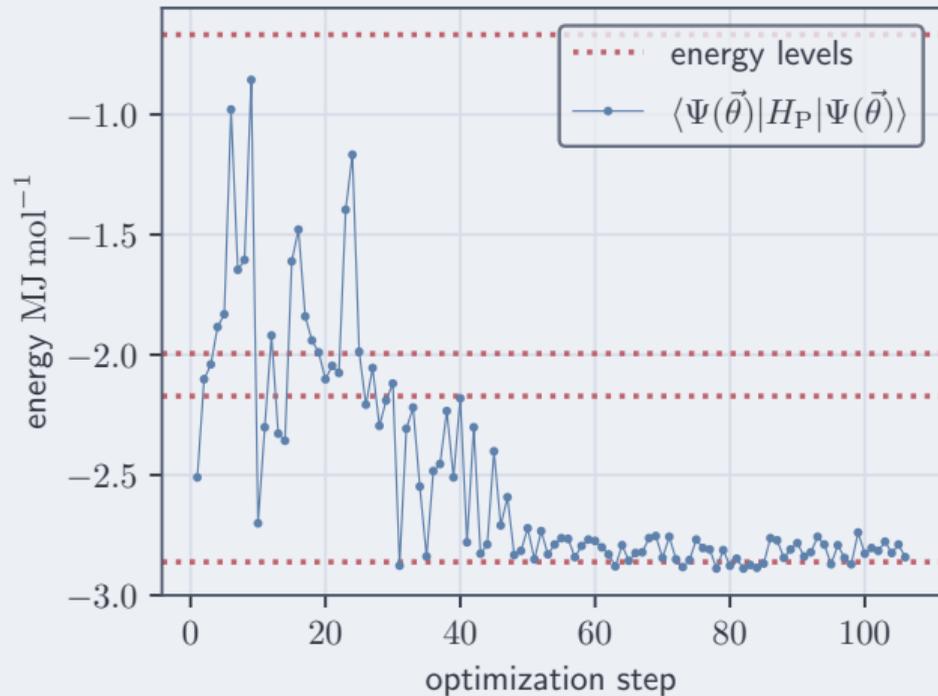
- trial state $|\Psi(\vec{\theta})\rangle$ is the UCC ansatz translated to Pauli matrices
- problem Hamiltonian H_P is the molecular Hamiltonian transformed to Pauli matrices
- optimization parameters are the UCC coefficients
- measurement of the observable can be achieved by clever measurement of Pauli strings

VQE: Example $\text{He} - \text{H}^+$



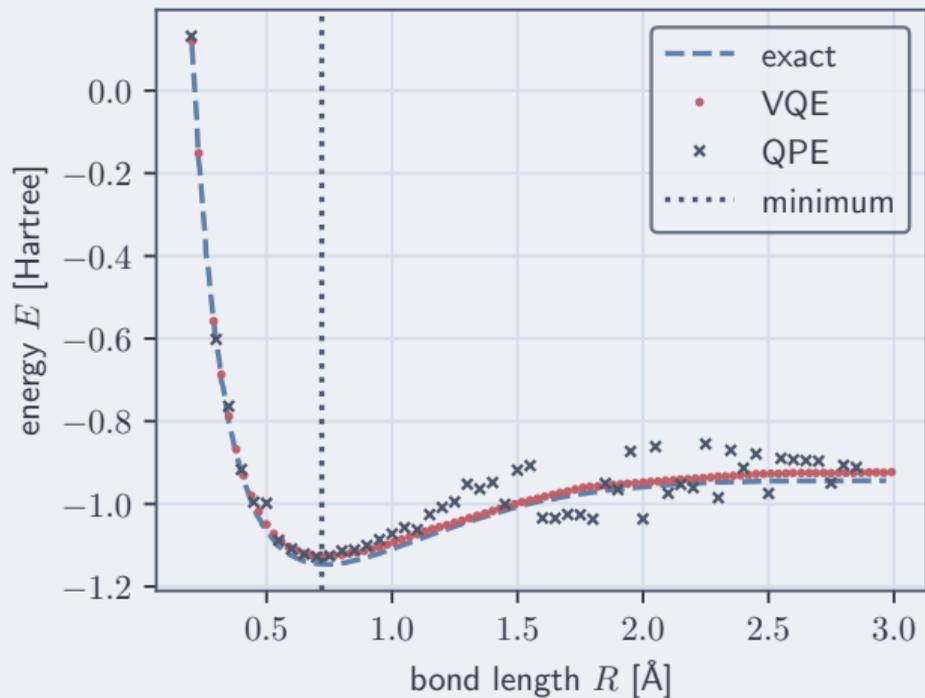
Data from Peruzzo et al. (2014), $\text{He} - \text{H}^+$ ground state energy using two photonic qubits

VQE: Example He – H⁺ II



- convergence of the Nelder-Mead (simplex) algorithm
- gradient-free (evaluations of the cost function require costly state preparations)

VQE: Example H_2



Data from O'Malley et al. (2016), H_2 ground state energy using two superconducting qubits

Quantum Phase Estimation (QPE): Principle

Previous to the VQE ansatz QPE was used for electronic structure problems on quantum computers.

- given a unitary operator U , QPE estimates the phase ω in $U|\Psi\rangle = e^{2\pi i\omega}|\Psi\rangle$ of an eigenstate $|\Psi\rangle$

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- the qubits of an additional register are iteratively entangled with the system register using controlled unitary gates
 - the system qubits act as control bits
 - the unitary is applied with different fixed phases $2^0, 2^1, 2^2, \dots$
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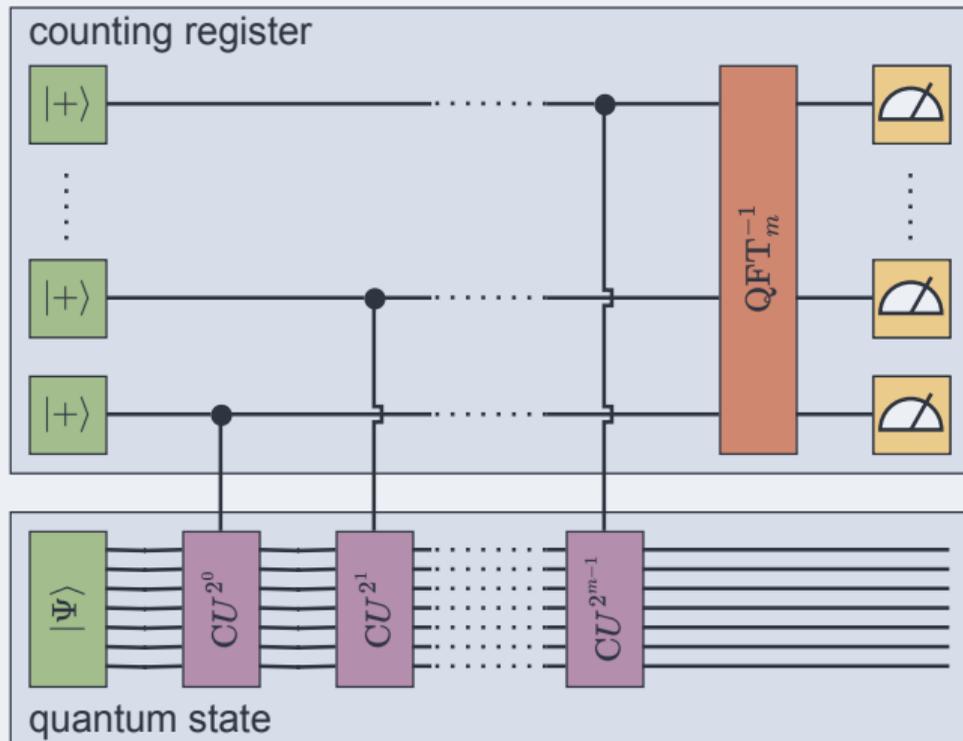
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 - the unitary is applied with different fixed phases $2^0, 2^1, 2^2, \dots$
 - the inverse quantum Fourier transformation is used to extract the phase from the ancillary register
- QPE reproduces ω with high probability within a margin ε using $\mathcal{O}(\log(1/\varepsilon))$ ancillary qubits
- requires $\mathcal{O}(1/\varepsilon)$ controlled unitary gates

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- the qubits of an additional register are iteratively entangled with the system register using controlled unitary gates
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 - the unitary is applied with different fixed phases $2^0, 2^1, 2^2, \dots$
 - the inverse quantum Fourier transformation is used to extract the phase from the ancillary register
- QPE reproduces ω with high probability within a margin ε using $\mathcal{O}(\log(1/\varepsilon))$ ancillary qubits
- requires $\mathcal{O}(1/\varepsilon)$ controlled unitary gates
- still requires a scheme to prepare $|\Psi\rangle$ close to the desired state

Quantum Phase Estimation (QPE): Sketch



Quantum Phase Estimation (QPE): Characteristics

Advantages:

- requires a single measurement
- in principle one could use QPE instead of QEE in VQE

Downsides:

- requires an extra qubit register
- the accuracy depends on the number of extra qubits m
- the number of required gates for the CU operations grows rapidly, quickly becomes infeasible due to limited decoherence times and gate fidelities
- CU gates are hard to implement
- if $2^m \omega$ is not an integer success probability is reduced from 1 to $\geq \frac{4}{\pi^2} \approx 0.405$

VQE: Classical Binary Optimization Problems

- optimization problem whose solutions are bit strings (N bits)

$$\mathbf{z} = z_1 z_2 \dots z_n \quad z_i \in \{0, 1\} \quad (31)$$

- How to generate trial states?

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$$C_1(\mathbf{z}) = z_1 \vee z_2, \quad C_2(\mathbf{z}) = z_2 \vee \bar{z}_3, \quad C_3(\mathbf{z}) = z_1 \vee \bar{z}_3 \vee z_4, \quad \dots \quad (33)$$

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- each clause typically depends only on a few bits
- $z_i \rightarrow \frac{1}{2}(1 - Z_i)$ maps cost function to an Ising Hamiltonian, can be treated with VQE ansatz

$$H_P = \sum_{i=1}^N h_i Z_i + \sum_{i,j=1}^N h_{ij} Z_i Z_j + \sum_{ijk=1}^N h_{ijk} Z_i Z_j Z_k + \dots \quad (33)$$

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Quantum Approximate Optimization Algorithm (QAOA)

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- many classical binary optimization problems are computationally hard
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- if $r \geq r^*$ for all problem instances: r^* characterizes the quality of the approximation algorithm

QAOA offers a heuristic ansatz to solve classical binary optimization problems encoded as Ising Hamiltonians

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and employs a mixing Hamiltonian

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A single integer parameter p (circuit depth) controls the trial state ansatz

$$|\Psi(\vec{\beta}, \vec{\gamma})\rangle = e^{-i\beta_p H_M} e^{-i\gamma_p H_P} \dots e^{-i\beta_1 H_M} e^{-i\gamma_1 H_P} |\Psi_0\rangle \quad (37)$$

that is determined by $2p$ parameters $0 \leq \gamma_i < 2\pi$ and $0 \leq \beta_i \leq \pi$.

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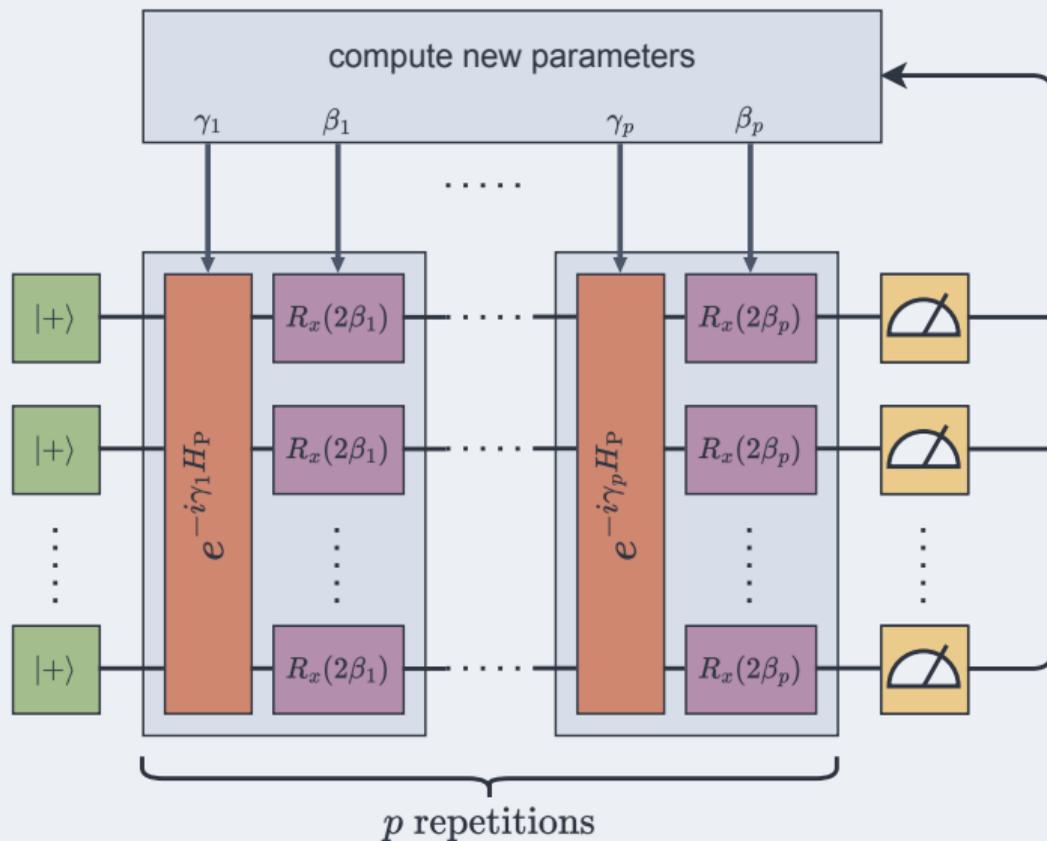
$$H_M = \sum_{i=1}^N X_i. \quad e^{-i\beta_j H_M} = \prod_{i=1}^N e^{-i\beta_j X_i} = \bigotimes_{i=1}^N R_x(2\beta_j) \quad (36)$$

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QAOA: Circuit Visualization



QAOA: Optimization Problem

- the parameters $\vec{\beta}$ and $\vec{\gamma}$ are optimized to maximize the expectation value

$$F_p = \langle \Psi(\vec{\beta}, \vec{\gamma}) | H_P | \Psi(\vec{\beta}, \vec{\gamma}) \rangle \quad (38)$$

- approximation ratio

$$r_p = \frac{F_p}{C_{\max}} \quad (39)$$

- adding a layer cannot worsen the result

$$F_{p+1} \geq F_p \quad (40)$$

- optimal limit

$$\lim_{p \rightarrow \infty} F_p = C_{\max} \quad (41)$$

approximates adiabatic quantum computing

QAOA: Motivation I

- recall adiabatic quantum computing:

$$H(t) = (1 - f(t))H_M + f(t)H_P \quad (42)$$

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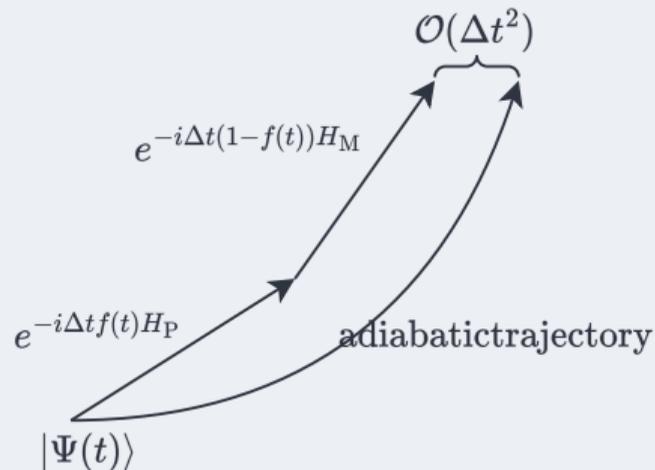
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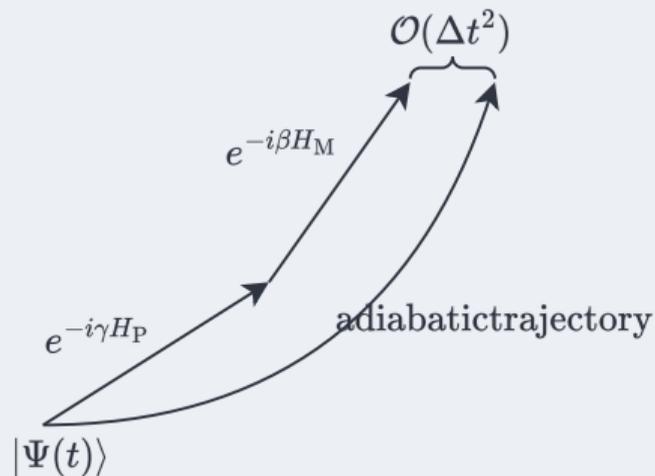
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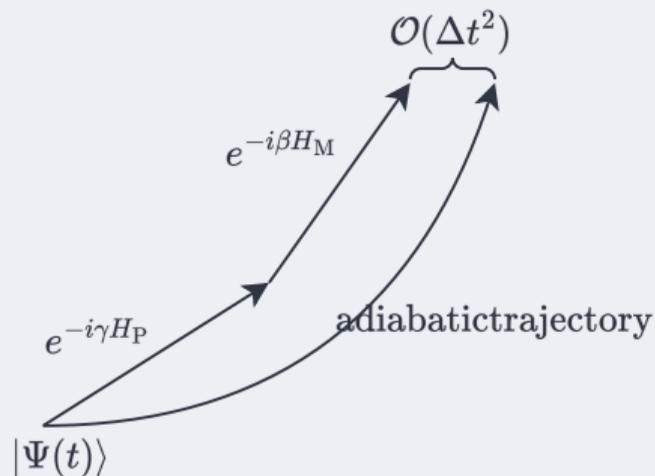
- define $\beta = (1 - f(t)) \Delta t$ and $\gamma = f(t) \Delta t$



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$$|\Psi(t + \Delta t)\rangle \approx \underbrace{\exp[-i\beta H_M] \exp[-i\gamma H_P]}_{\text{one QAOA layer}} |\Psi(t)\rangle + \mathcal{O}(\Delta t^2) \quad (46)$$

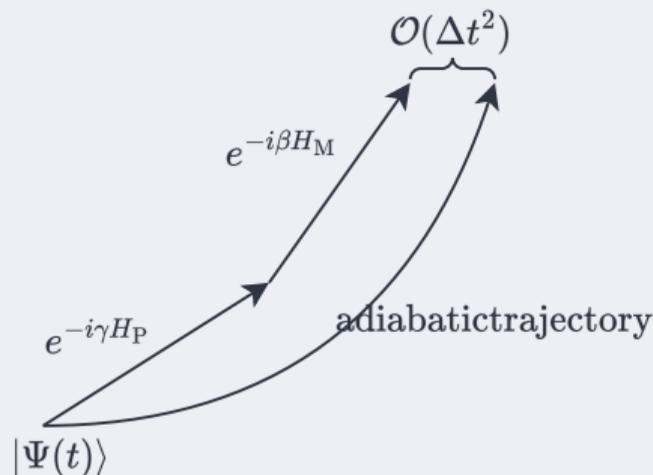
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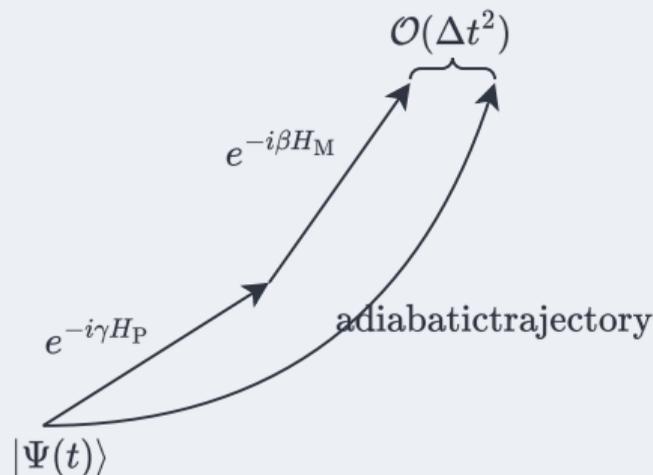
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- however QAOA optimizes β and γ independently
- $p \rightarrow \infty$ (infinitely many layers) reproduces adiabatic trajectory



QAOA: Applications

- maximum likelihood channel decoding (communications)
- community detection (social, neural and biological networks)
- portfolio optimization (finance industry)
- exact cover problem (tiling problems and aircraft flight gate assignment)
- maximum independent set (radio networks and genetic engineering)
- MaxCut (integrated circuit design, statistical physics, data clustering)
- linear algebra (fundamental for many applications)
- traveling salesperson (logistics, traveling)
- graph coloring (scheduling, compiler optimization)
- correlation clustering (data science)
- other satisfiability problems, ...

QAOA: Advantages

- well controllable
 - single integer parameter determines the circuit structure
 - increasing p can not worsen solution (monotonous)
- it has been shown that the parametric square-pulse ansatz (“bang-bang”) is optimal given a fixed quantum computation time budget
- QAOA can learn to exploit diabatic transitions to overcome small energy gap limitations
- gate imperfections can be mitigated by parameter tuning
- circuits are relatively shallow
- nice implementation on Rydberg platform (driving with global pulses)

Max-Cut

Max-Cut: Problem Statement

Partition the vertices of a graph in two sets, such that the sum of edge weights between the sets becomes maximal.

$$C(\mathbf{z}) = \sum_{i,j=1}^N w_{ij} z_i (1 - z_j) \quad (47)$$

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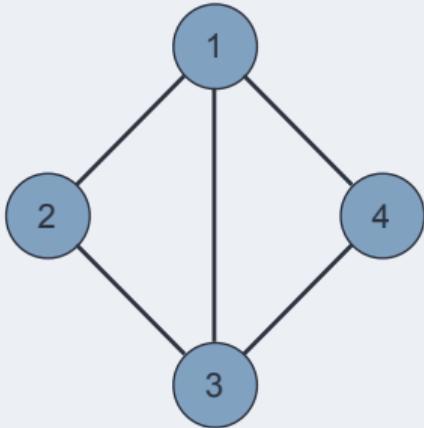
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- stochastic solvers for $N = \mathcal{O}(10^2)$ (no solution guaranteed)
- APX-hard: every polynomial-time approximation scheme has an approximation ratio guarantee $r^* < 1$
- solving the problem for generic graphs with $r^* \geq 16/17 \approx 0.9412$ is NP-hard
- the best classical algorithm can only guarantee $r^* \gtrsim 0.87856$

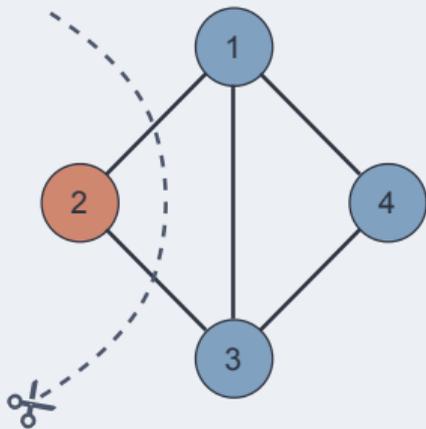
Max-Cut: Example

Here: equal weights $w_{ij} = 1$



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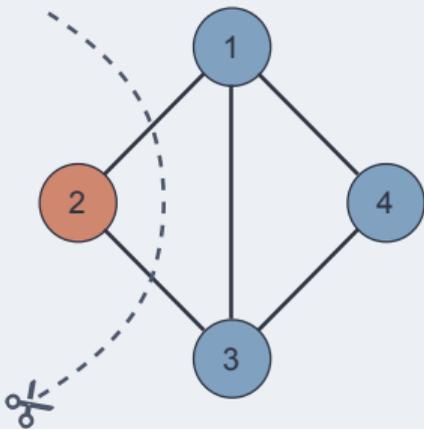
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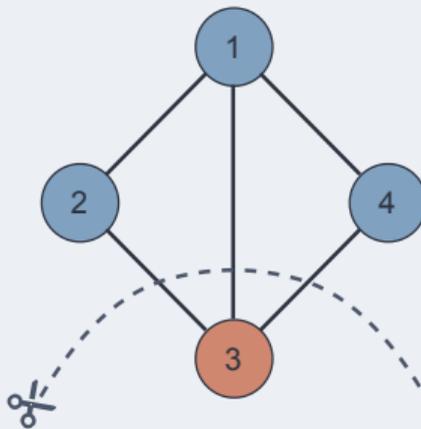
two edges

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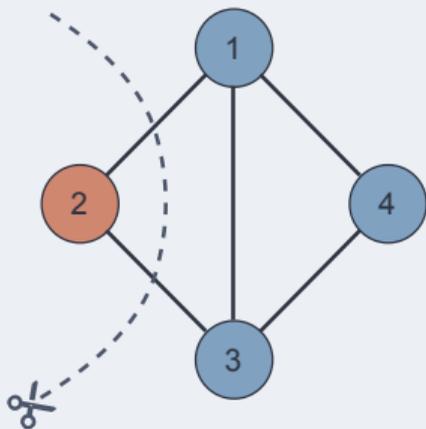
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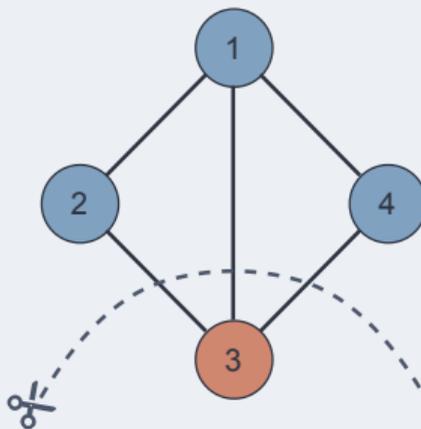
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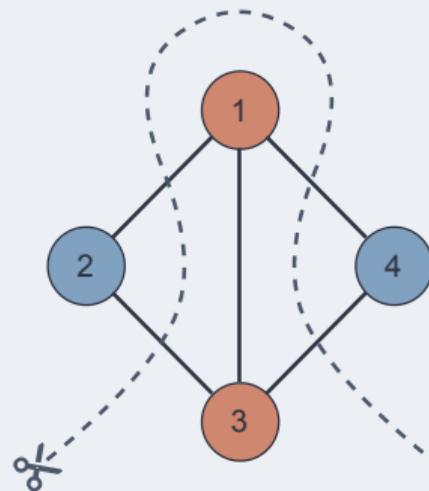
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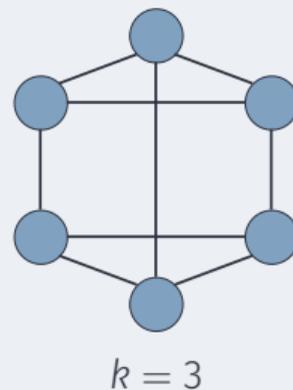
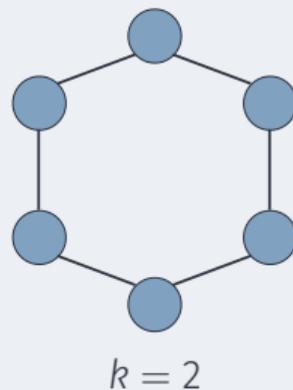
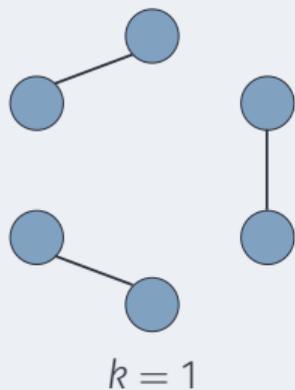
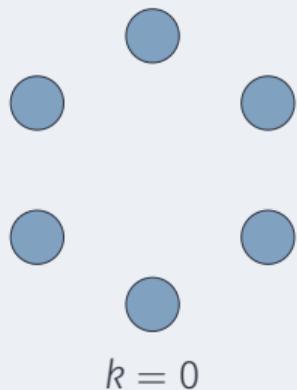
four edges

Max-Cut: k -Regular Graphs

- often: only consider k -regular graphs (usually $k = 3$), i.e. graphs where each vertex has exactly 3 edges to other vertices.

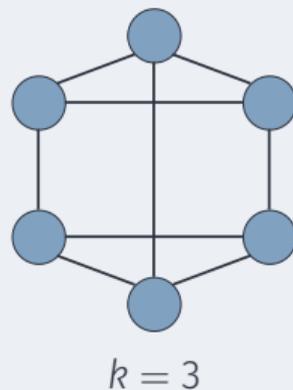
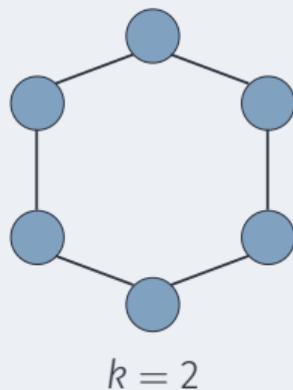
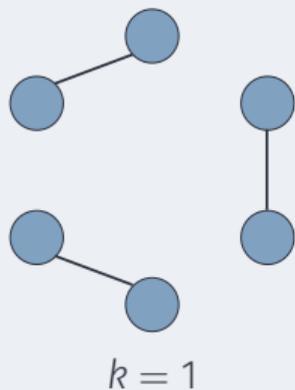
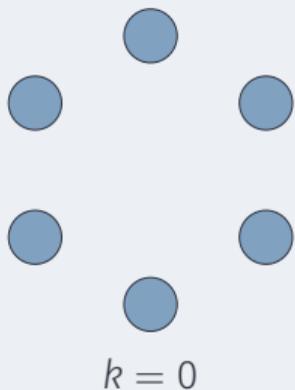
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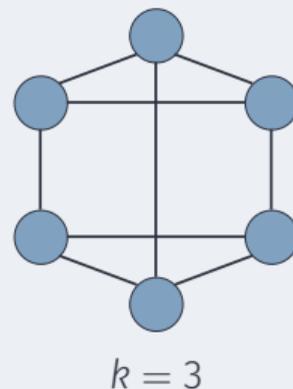
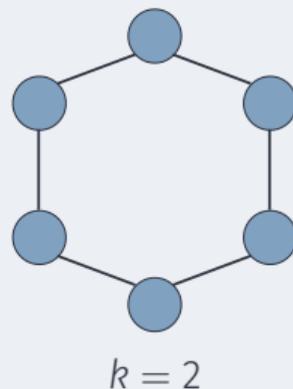
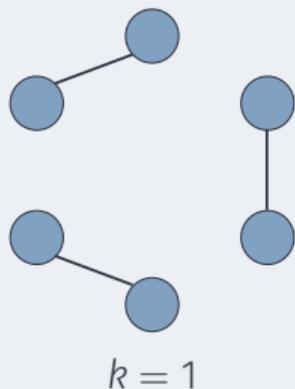
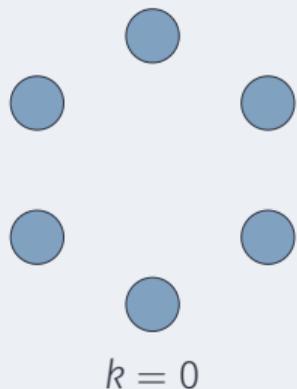
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- solving with $r^* \geq 331/332 \approx 0.9970$ is NP-hard
- best classical algorithm guarantees $r \gtrsim 0.9326$



Max-Cut: Application of QAOA

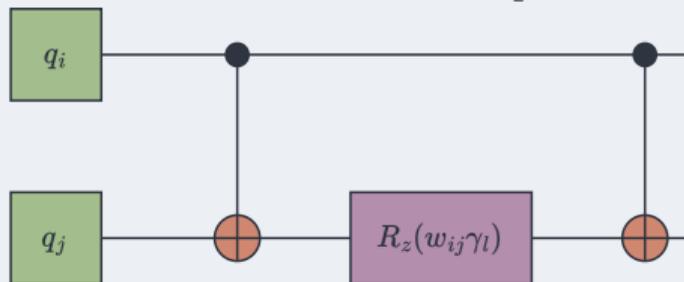
- many studies investigating this problem (including original QAOA paper)
- the problem can be encoded as the Ising Hamiltonian

$$H_P = \frac{1}{2} \sum_{\langle i,j \rangle} w_{ij} (1 - Z_i Z_j) \quad (48)$$

- problem instance specific constant

$$C = \frac{1}{2} \sum_{\langle i,j \rangle} w_{ij} \quad (49)$$

- sub-circuit for each edge of the graph ($w_{ij} = 1$): $\exp \left[-\frac{i}{2} w_{ij} \gamma_l Z_i Z_j \right]$



Max-Cut: QAOA Results

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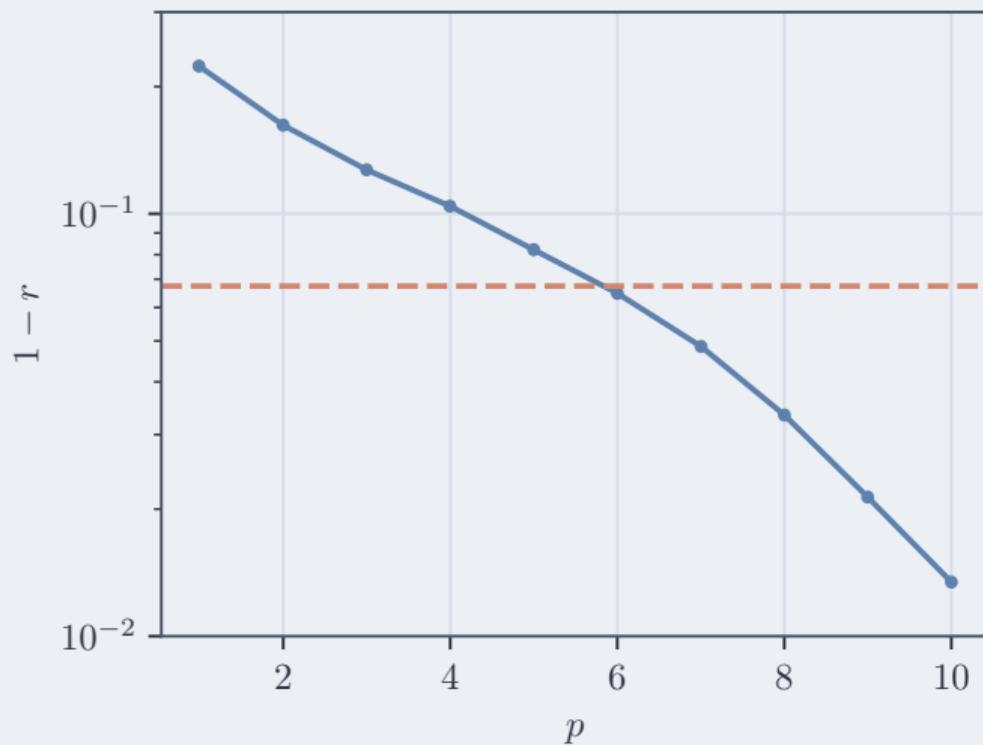
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Zhou et al. (2020):

- robustness against small energy gaps
- exploit patterns in parameters

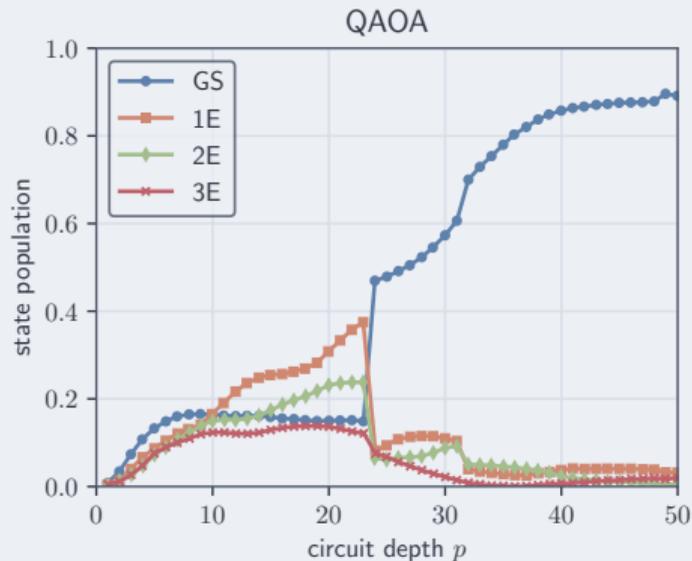
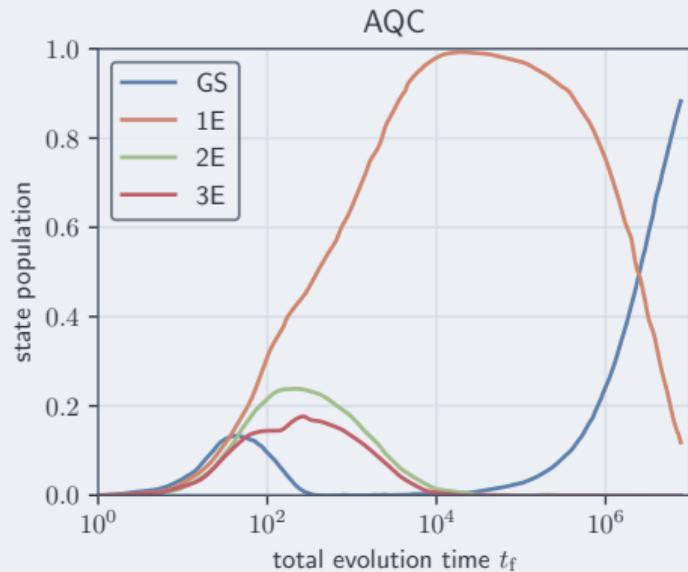
Max-Cut: Approximation Ratio for Different p

Performance on weighted 4-regular graphs with 16 vertices, data from Zhou et al. (2020)



Max-Cut: Small Energy Gaps

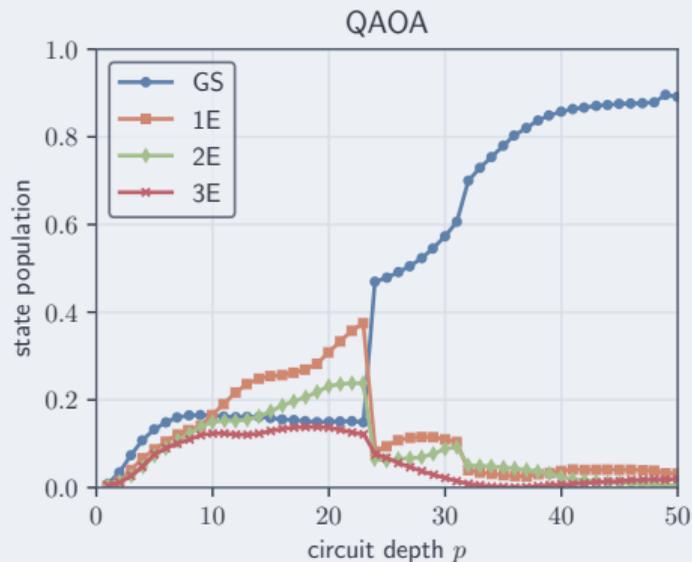
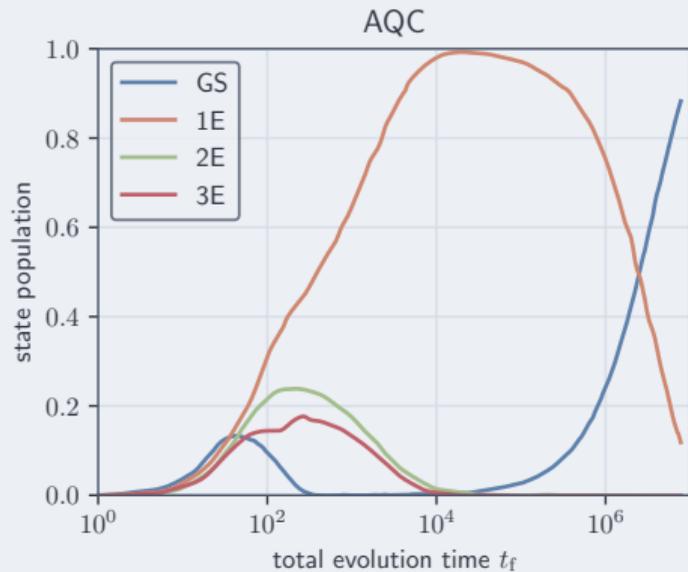
AQC requires extremely long evolution times when using a simple trajectory



Data from Zhou et al. (2020), 3-regular graph with $N = 14$ nodes and a small minimum spectral gap $\Delta E < 10^{-3}$

Max-Cut: Small Energy Gaps

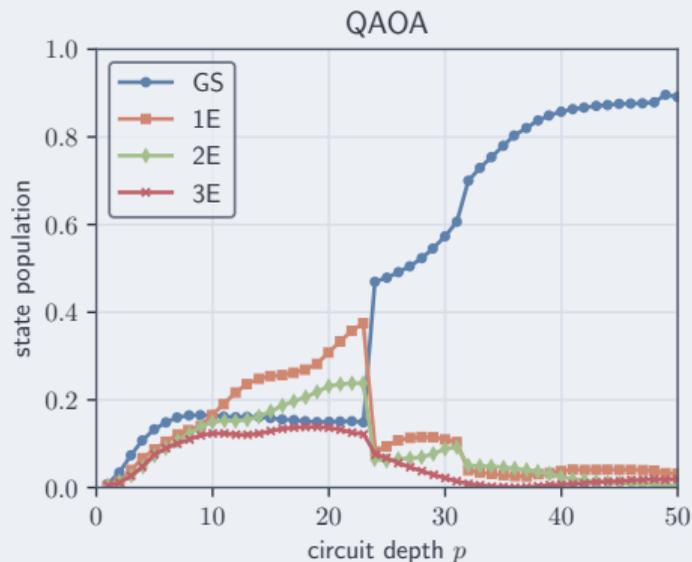
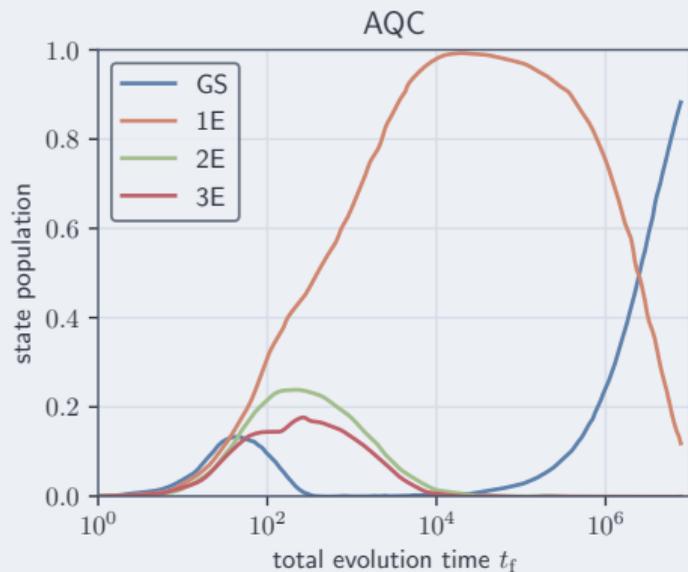
QAOA is reasonably robust against the small energy gap



Data from Zhou et al. (2020), 3-regular graph with $N = 14$ nodes and a small minimum spectral gap $\Delta E < 10^{-3}$

Max-Cut: Small Energy Gaps

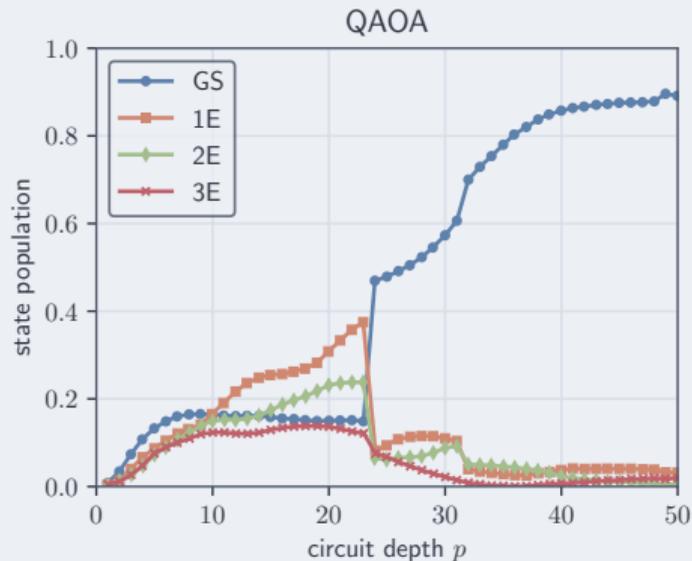
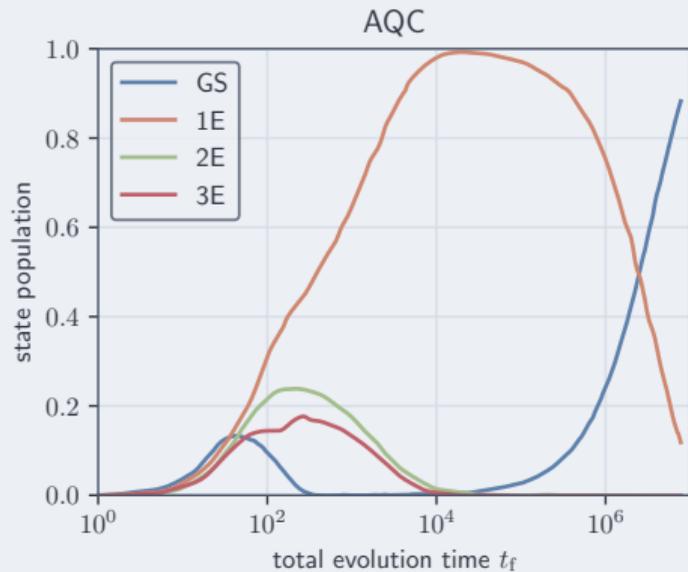
$p \gtrsim 24$ significantly improves the result



Data from Zhou et al. (2020), 3-regular graph with $N = 14$ nodes and a small minimum spectral gap $\Delta E < 10^{-3}$

Max-Cut: Small Energy Gaps

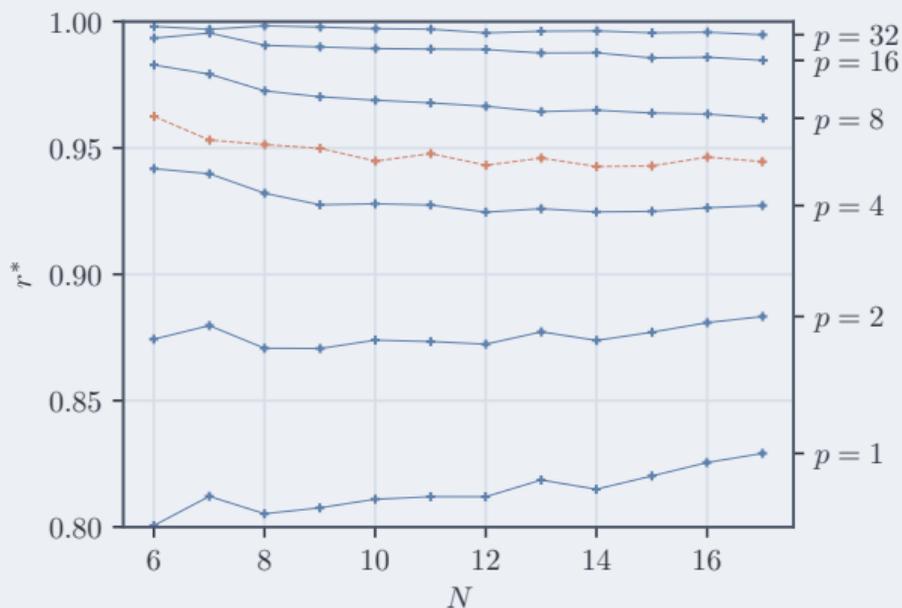
QAOA result can be used to craft an optimized trajectory for AQC



Data from Zhou et al. (2020), 3-regular graph with $N = 14$ nodes and a small minimum spectral gap $\Delta E < 10^{-3}$

Max-Cut: Random Graphs (Erdős–Rényi)

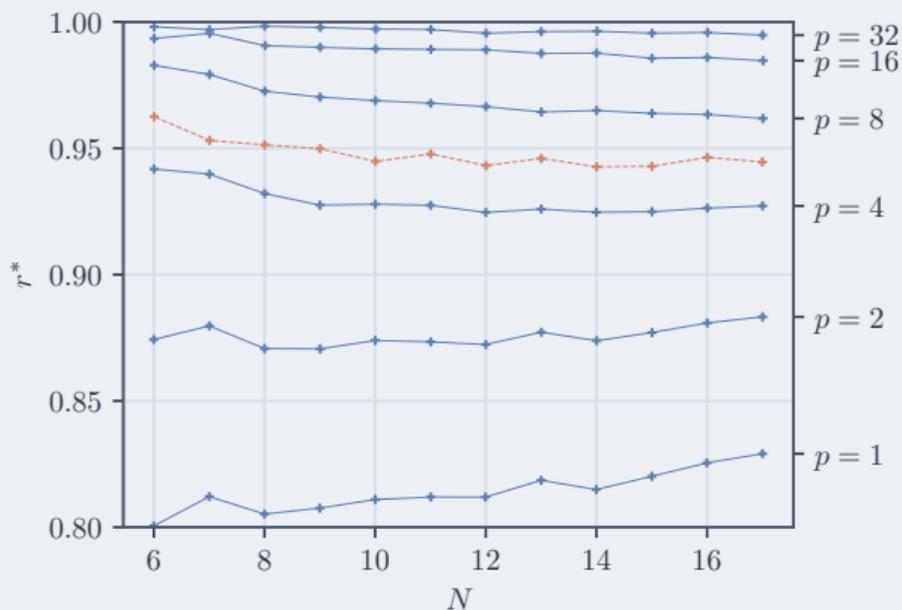
QAOA can also perform well on random graphs



orange: best classical algorithm, data from Crooks (2018)

Max-Cut: Random Graphs (Erdős–Rényi)

Here $p \geq 8$ achieves quantum advantage (note system size however!)



orange: best classical algorithm, data from Crooks (2018)

Conclusions & Outlook

Conclusions:

- VQE has been used to accurately reproduce the energies of simple molecules
- with current hardware it produces more accurate results than QPE

Conclusions & Outlook: VQE

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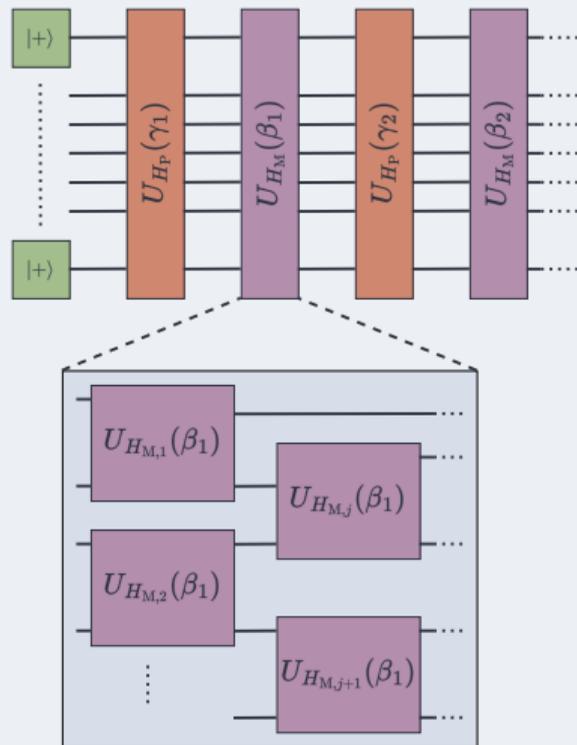
Outlook:

- treatment of larger molecules:
 - largest so far 6 qubits for BeH_2 (heuristic ansatz instead of UCC)
 - far from any actual quantum advantage
- deeper analysis of convergence and local optima

Outlook: QAOA

Outlook:

- better understanding for $p > 1$
- optimizing mixer Hamiltonian
- improved strategies for optimizing $\vec{\gamma}/\vec{\beta}$
- tackle barren plateaus
- demonstrate quantum advantage
 - surpass classical approximation ratio on system sizes that cannot be treated with exact algorithms
 - actual experimental implementations
 - “soon”: QAOA for $N = 400$ vertices with $p \gtrsim 25$ (Rydberg atoms in optical tweezer arrays)



Select References: VQE

- Original work, $\text{He} - \text{H}^+$ using two photonic qubits
Alberto Peruzzo et al. **“A Variational Eigenvalue Solver on a Photonic Quantum Processor”**. en. In: *Nature Communications* 5.1 (July 2014), p. 4213. ISSN: 2041-1723. DOI: 10/f6df9g
- H_2 using two superconducting qubits
P. J. J. O’Malley et al. **“Scalable Quantum Simulation of Molecular Energies”**. In: *Phys. Rev. X* 6.3 (July 2016), p. 031007. DOI: 10/f8wcgg
- theoretical discussion of the method
Jarrod R. McClean et al. **“The Theory of Variational Hybrid Quantum-Classical Algorithms”**. en. In: *New J. Phys.* 18.2 (Feb. 2016), p. 023023. ISSN: 1367-2630. DOI: 10/gdcc5s

Select References: QAOA

- Original work, Max-Cut with QAOA
Edward Farhi, Jeffrey Goldstone, and Sam Gutmann. **“A Quantum Approximate Optimization Algorithm”**. In: *arXiv:1411.4028 [quant-ph]* (Nov. 2014). arXiv: 1411.4028 [quant-ph]
- In-depth analysis, comparison to AQC
Leo Zhou et al. **“Quantum Approximate Optimization Algorithm: Performance, Mechanism, and Implementation on Near-Term Devices”**. In: *Phys. Rev. X* 10.2 (June 2020), p. 021067. DOI: 10/gg4nk2
- Random graphs
Gavin E. Crooks. **“Performance of the Quantum Approximate Optimization Algorithm on the Maximum Cut Problem”**. In: *arXiv:1811.08419 [quant-ph]* (Nov. 2018). arXiv: 1811.08419 [quant-ph]

Thank you for your attention!